

# Stress Relaxation Due to Dislocation Formation in Orthorhombic Ga<sub>2</sub>O<sub>3</sub> Films Grown on Al<sub>2</sub>O<sub>3</sub> Substrates

A.M. Smirnov<sup>1,\*</sup>, A.Yu. Ivanov<sup>1</sup>, A.V. Kremleva<sup>1</sup>, Sh.Sh. Sharofidinov<sup>2</sup>, A.E. Romanov<sup>1,3</sup>

<sup>1</sup> Institute of Advanced Data Transfer Systems, ITMO University, Kronverkskiy pr. 49, St. Petersburg 197101, Russia

<sup>2</sup> Sector of Solid State Electronics, Ioffe Physical-Technical Institute, Russian Academy of Sciences, Polytechnicheskaya 26, St. Petersburg, 194021, Russia

<sup>3</sup> Sector of Solid State Theory, Ioffe Physical-Technical Institute, Russian Academy of Sciences, Polytechnicheskaya 26, St. Petersburg, 194021, Russia

## Article history

Received September 20, 2022

Accepted September 21, 2022

Available online September 25, 2022

## Abstract

We analyze the preference of various types of misfit dislocation (MD) formation in film/substrate  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and  $\kappa$ -(Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub> heterostructures. We consider two possibilities for variation in films growth orientation (defined by inclination angle  $\vartheta$ ) for these heterostructures with inclination axes about either [100] or [010] crystallographic directions. We study dependences of the critical film thickness for MD formation on the inclination angle  $\vartheta$  for heterostructures under consideration. We find the presence of two special orientations ( $\vartheta \sim 26^\circ$  for [100] heterostructure,  $\vartheta \sim 28^\circ$  for [010] heterostructure, and  $\vartheta = 90^\circ$  for both inclination types) of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> heterostructures, for which the formation of MDs is energetically unfavorable. We show that formation of pure edge MDs is easier for [010]  $\kappa$ -(Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub> heterostructures than for [100] heterostructures, and it is vice versa for mixed MDs in these heterostructures.

**Keywords:** Film/substrate heterostructures; Gallium oxide; Growth orientation; Misfit dislocations

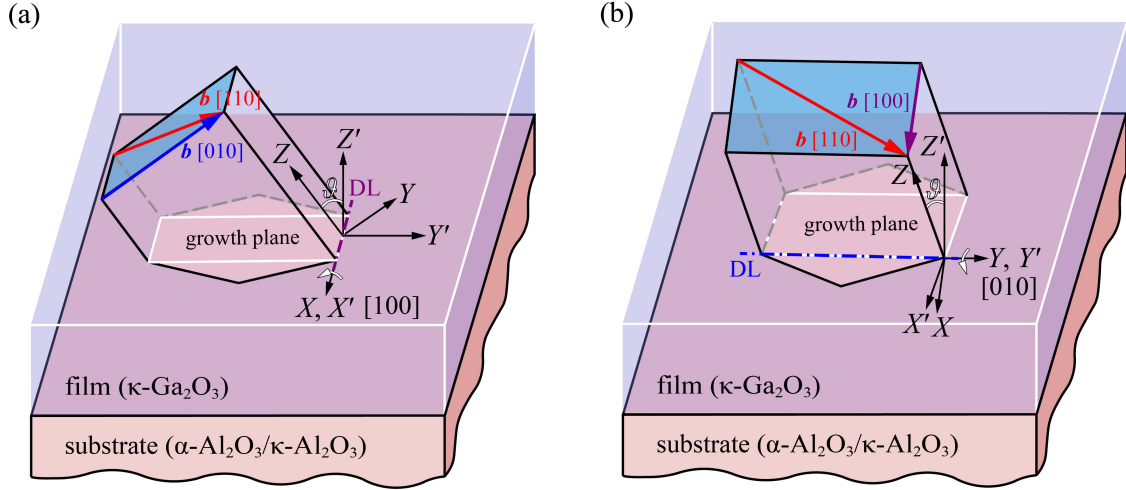
## 1. INTRODUCTION

Gallium oxide (Ga<sub>2</sub>O<sub>3</sub>), one of the ultra-wide bandgap semiconductors, has attracted much attention in recent years due to great prospects in power electronic devices including those for space and nuclear industries, photoelectric converters for the ultraviolet region, high-power radio frequency devices, gas sensors, etc. [1–3]. An important feature of Ga<sub>2</sub>O<sub>3</sub> is the possibility of its existence in various polymorphic forms: stable monoclinic  $\beta$ -phase, and metastable  $\alpha$ -,  $\gamma$ -,  $\kappa$ -,  $\varepsilon$ - and  $\delta$ -phases with their specific characteristics and benefits [4–8]. The metastable hexagonal  $\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub> is of particular interest due to its favorable growth features and great compatibility with common substrates and the possibility of polarization engineering endowed by its polar nature [9]. Early reports investigated a hexagonal  $P6_3mc$  structure for  $\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub>, while recent investigations identified its orthorhombic structure in the  $Pna2_1$  symmetry group, denoted as  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> [10,11]. The pseudo-hexagonal structure of  $\varepsilon$ -phase consists of 120°

rotational orthorhombic  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> domains. However, until now, some works indicate the possibility of the existence of pure  $\varepsilon$ -phase [12]. Despite the ongoing debate about the real structure of the considered phase, most scientific groups tend to consider the phase as composed of 3-fold rotational orthorhombic domains with disordered domain boundaries [10,11,13,14]; the single-domain orthorhombic structure of this phase has been recently demonstrated [3,15]. Therefore, in this work, we will consider this phase as  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> with the orthorhombic crystal structure  $Pna2_1$ .

The heterostructures based on  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> are successfully used in solar blind photodetectors [16], powerful Schottky diodes [17], and in high-frequency applications [18]. Unfortunately, the characteristics of these devices are limited by the crystal quality of the  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> film/layer/heterostructures, which are characterized by a high dislocation density that ranges from  $5.2 \times 10^7$  to  $1.4 \times 10^{13}$  cm<sup>-2</sup> [3,19]. Such a wide scatter in the value of the dislocation density can be due to the insufficient

\* Corresponding author: A.M. Smirnov, e-mail: [smirnov.mech@gmail.com](mailto:smirnov.mech@gmail.com)



**Fig. 1.** Schematics of the film/substrate heterostructure with selected coordinate systems that are used in the calculations. Frame (a) shows heterostructure with inclination axis about [100] crystallographic direction for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/κ-Al<sub>2</sub>O<sub>3</sub> (i.e., about  $[2\bar{1}\bar{1}0]$  for  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>); frame (b) shows heterostructure with inclination axis about [010] crystallographic direction for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/κ-Al<sub>2</sub>O<sub>3</sub> (i.e., about  $[\bar{1}2\bar{1}0]$  for  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>).  $\vartheta$  is the angle between the  $c$ -axis ( $Z$ -axis) of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> and the normal one to the film plane ( $Z'$ -axis). Blue surfaces show (001) crystallographic plane of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>. Red, blue and purple arrows show the direction of Burgers vector  $\mathbf{b}$  along [110], [010] and [100] crystallographic directions, respectively. Dash-dotted lines indicate dislocation lines (DLs).

knowledge on the optimal parameters of the heterostructure, including substrate material, crystallographic orientation of the substrate, film thickness, etc. in order to grow  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> with the minimum possible number of defects. In this work, we propose a theoretical model that allows to find the critical conditions of the system for the formation of misfit dislocations (MDs) in  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/α-Al<sub>2</sub>O<sub>3</sub> and  $\kappa$ -(Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub>/κ-Al<sub>2</sub>O<sub>3</sub> heterostructures.

## 2. MODEL

Similarly to the approach described in Refs. [20–25], we consider a film/substrate  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub> heterostructure with thin film and MDs formed at the interface (Fig. 1). Since the film consists of the  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>, it is possible to assume the following variation in heterostructure growth orientations with minimal mismatch between  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> film and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> or  $\kappa$ -Al<sub>2</sub>O<sub>3</sub> substrate: (i) heterostructure with inclination axes about  $\langle 100 \rangle$ -type crystallographic directions for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/κ-Al<sub>2</sub>O<sub>3</sub> (i.e. about  $\langle 2\bar{1}\bar{1}0 \rangle$  for  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>), see Fig. 1a, and (ii) heterostructure with inclination axes about  $\langle 010 \rangle$ -type crystallographic directions for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/κ-Al<sub>2</sub>O<sub>3</sub> (i.e. about  $\langle \bar{1}100 \rangle$  for  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>), see Fig. 1b. The inclination angle is given as  $\vartheta$  in Fig. 1.

In general, the MDs formed at the heterointerface as a result of stress relaxation process are mixed dislocations, which have both edge and screw Burgers vector components. The Burgers vectors (colored arrows) and dislocation lines (DLs; colored dash-dotted lines) of such MDs are shown in Fig. 1. Red, blue and purple arrows correspond to Burgers vector along [111], [010] and [100]

crystallographic directions, respectively; DLs are parallel to inclination axes. In this work, we consider only mixed and pure edge MDs; pure screw MDs do not contribute to the misfit stresses relaxation in the case of a biaxial stress-strain state under consideration [26].

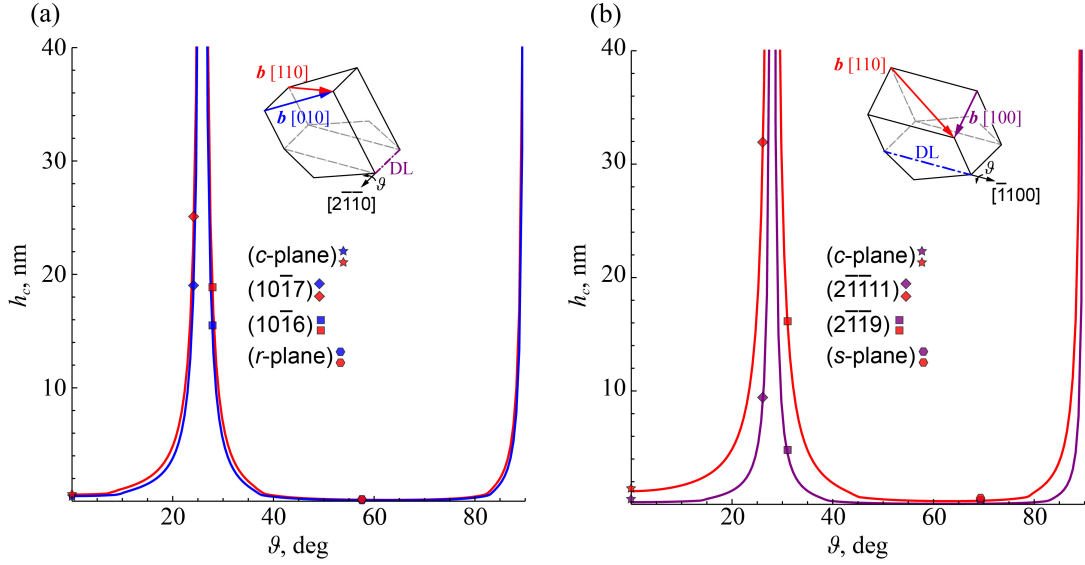
Our model, which describes the critical conditions for the MDs formation in  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/α-Al<sub>2</sub>O<sub>3</sub> and  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/κ-Al<sub>2</sub>O<sub>3</sub> heterostructures, is based on the Matthews-Blakeslee (M-B) approach [27,28]. For the elastically semi-isotropic case, the M-B approach gives the simplified expression for finding the critical thickness  $h_c$  for MD formation [20]:

$$h_c = \frac{b_{\parallel}^2 + b_{\perp}^2 + (1-\nu)b_s^2}{\varepsilon_m(1+\nu)8\pi b_{\parallel}} \ln\left(\frac{2h_c}{r_0}\right), \quad (1)$$

where  $\varepsilon_m$  is the component of the lattice mismatch released by MDs;  $\nu$  is Poisson ratio ( $\nu = 0.3$  in our calculations);  $b_{\parallel}$  and  $b_{\perp}$  are the projections of the edge component of the dislocation Burgers vector parallel and perpendicular heterointerface, respectively, and  $b_s$  is the magnitude of the screw component of the dislocation Burgers vector;  $r_0$  is the dislocation core radius parameter. The magnitude of the Burgers vector  $\mathbf{b}$  is related to the components:  $b^2 = b_{\parallel}^2 + b_{\perp}^2 + b_s^2$ .

The anisotropy of the film and substrate materials is considered in the parameter  $\varepsilon_m$ , which takes the following form:

$$\varepsilon_m = \frac{a_{\alpha\text{-Al}_2\text{O}_3} \sqrt{3} c_{\alpha\text{-Al}_2\text{O}_3} - \sqrt{(b_{\kappa\text{-Ga}_2\text{O}_3} c_{\alpha\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + 3(a_{\alpha\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}{\sqrt{(b_{\kappa\text{-Ga}_2\text{O}_3} c_{\alpha\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + 3(a_{\alpha\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}} \quad (2)$$



**Fig. 2.** Dependences of the critical thickness  $h_c$  on the inclination angle  $\vartheta$  for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> heterostructures. Frames (a) and (b) show dependences for heterostructures with inclination axes about  $[2\bar{1}\bar{1}0]$  and  $[\bar{1}\bar{1}00]$  crystallographic direction of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> substrate, respectively. Red, blue and purple curves describe the cases of MDs with Burgers vector  $\mathbf{b}$  along  $[110]$ ,  $[010]$  and  $[100]$  crystallographic directions, respectively. Bullets indicate  $h_c$  values for various growth planes of the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> substrate.

for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> heterostructure with inclination axis about  $[100]$  crystallographic direction of the film or  $[2\bar{1}\bar{1}0]$  crystallographic direction of the substrate;

$$\varepsilon_m = \frac{a_{\alpha\text{-Al}_2\text{O}_3} c_{\alpha\text{-Al}_2\text{O}_3} - \sqrt{(a_{\kappa\text{-Ga}_2\text{O}_3} c_{\alpha\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + (a_{\alpha\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}{\sqrt{(a_{\kappa\text{-Ga}_2\text{O}_3} c_{\alpha\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + (a_{\alpha\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}, \quad (3)$$

for heterostructure  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> with inclination axis about  $[010]$  crystallographic direction of the film or  $[\bar{1}\bar{1}00]$  crystallographic direction of the substrate;

$$\varepsilon_m = \frac{b_{\kappa\text{-Al}_2\text{O}_3} c_{\kappa\text{-Al}_2\text{O}_3} - \sqrt{(b_{\kappa\text{-Ga}_2\text{O}_3} c_{\kappa\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + (b_{\kappa\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}{\sqrt{(b_{\kappa\text{-Ga}_2\text{O}_3} c_{\kappa\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + 3(b_{\kappa\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}, \quad (4)$$

for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub> heterostructure with inclination axis about  $[100]$  crystallographic direction;

$$\varepsilon_m = \frac{a_{\kappa\text{-Al}_2\text{O}_3} c_{\kappa\text{-Al}_2\text{O}_3} - \sqrt{(a_{\kappa\text{-Ga}_2\text{O}_3} c_{\kappa\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + (a_{\kappa\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}{\sqrt{(a_{\kappa\text{-Ga}_2\text{O}_3} c_{\kappa\text{-Al}_2\text{O}_3})^2 \cos^2 \vartheta + (a_{\kappa\text{-Al}_2\text{O}_3} c_{\kappa\text{-Ga}_2\text{O}_3})^2 \sin^2 \vartheta}}, \quad (5)$$

for heterostructure  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub> with inclination axis about  $[010]$  crystallographic direction  $a_{\kappa\text{-Ga}_2\text{O}_3}$ ,  $b_{\kappa\text{-Ga}_2\text{O}_3}$ ,  $c_{\kappa\text{-Ga}_2\text{O}_3}$ ,  $a_{\alpha\text{-Al}_2\text{O}_3}$ ,  $c_{\alpha\text{-Al}_2\text{O}_3}$ ,  $a_{\kappa\text{-Al}_2\text{O}_3}$ ,  $b_{\kappa\text{-Al}_2\text{O}_3}$ ,  $c_{\kappa\text{-Al}_2\text{O}_3}$  are the crystal lattice parameters of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>,  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and  $\kappa$ -Al<sub>2</sub>O<sub>3</sub>; their values are given in Table 1.

The possible orientations of the Burgers vector are chosen based on the crystal structure of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> such as parallel to  $[100]$ ,  $[010]$  and  $[110]$  crystallographic directions. In the case of the heterostructures with inclination axis  $[100]$ , the projections of the Burgers vector have the form:

$$\mathbf{b} [100]: b_s = a_{\kappa\text{-Ga}_2\text{O}_3}, \quad b_{\parallel} = 0, \quad b_{\perp} = 0,$$

$$\mathbf{b} [010]: b_s = 0, \quad b_{\parallel} = b_{\kappa\text{-Ga}_2\text{O}_3} \cos \vartheta, \quad b_{\perp} = b_{\kappa\text{-Ga}_2\text{O}_3} \sin \vartheta,$$

$$\mathbf{b} [110]: b_s = a_{\kappa\text{-Ga}_2\text{O}_3}, \quad b_{\parallel} = b_{\kappa\text{-Ga}_2\text{O}_3} \cos \vartheta, \\ b_{\perp} = b_{\kappa\text{-Ga}_2\text{O}_3} \sin \vartheta, \quad (6)$$

and in the case of the heterostructures with inclination axis  $[010]$ , the projections of the Burgers vector have the form:

$$\mathbf{b} [100]: b_s = 0, \quad b_{\parallel} = a_{\kappa\text{-Ga}_2\text{O}_3} \cos \vartheta, \quad b_{\perp} = -a_{\kappa\text{-Ga}_2\text{O}_3} \sin \vartheta,$$

$$\mathbf{b} [010]: b_s = b_{\kappa\text{-Ga}_2\text{O}_3}, \quad b_{\parallel} = 0, \quad b_{\perp} = 0,$$

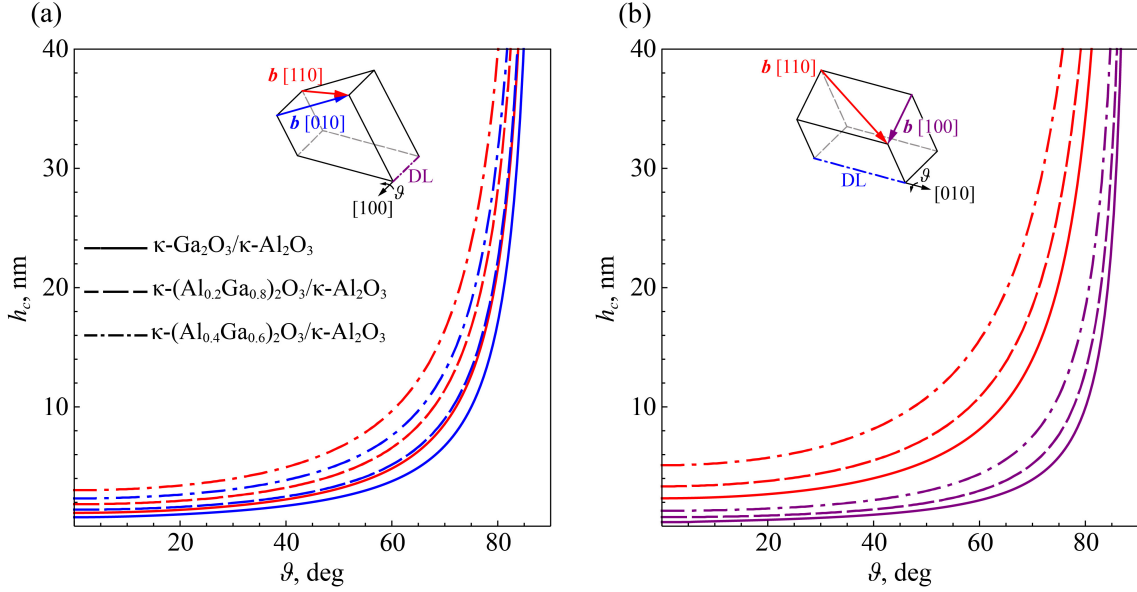
$$\mathbf{b} [110]: b_s = b_{\kappa\text{-Ga}_2\text{O}_3}, \quad b_{\parallel} = a_{\kappa\text{-Ga}_2\text{O}_3} \cos \vartheta, \\ b_{\perp} = -a_{\kappa\text{-Ga}_2\text{O}_3} \sin \vartheta. \quad (7)$$

### 3. RESULTS AND DISCUSSION

Figure 2 presents the calculated dependences of the critical film thickness  $h_c$  for the formation of various types of MDs on the inclination angle  $\vartheta$  for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> heterostructures. A feature of these dependences is the

**Table 1.** Crystal lattice parameters of materials under consideration.

Lattice parameters (Å)	$\kappa$ -Ga <sub>2</sub> O <sub>3</sub> [29] $Pna2_1$	$\alpha$ -Al <sub>2</sub> O <sub>3</sub> [30] $R\bar{3}c:H$	$\kappa$ -Al <sub>2</sub> O <sub>3</sub> [29] $Pna2_1$
$a$	5.078	4.761	4.848
$b$	8.706	4.761	8.323
$c$	9.308	12.996	8.944



**Fig. 3.** Dependences of the critical thickness  $h_c$  on the inclination angle  $\vartheta$  for  $\kappa$ -( $\text{Al}_x\text{Ga}_{1-x}$ ) $_2\text{O}_3/\kappa$ - $\text{Al}_2\text{O}_3$  heterostructures. Frames (a) and (b) show dependences for heterostructures with inclination axes about  $[100]$  and  $[010]$  crystallographic direction of  $\kappa$ - $\text{Al}_2\text{O}_3$  substrate, respectively. Solid, dashed and dash-dotted curves correspond to the value of Al composition  $x = 0, 0.2,$  and  $0.4$  respectively.

presence of two special orientations of heterostructures in which the formation of MDs is energetically unfavorable. First special case takes place at  $\vartheta \sim 26^\circ$  (for heterostructure with inclination axis about  $[2\bar{1}\bar{1}0]$ ) and  $\vartheta \sim 28^\circ$  (for heterostructure with inclination axis about  $[\bar{1}100]$ ); formation of MDs is impossible in such oriented heterostructures due to vanishing of  $\varepsilon_m$  because of the values of  $\kappa$ - $\text{Ga}_2\text{O}_3$  and  $\alpha$ - $\text{Al}_2\text{O}_3$  crystal lattice parameters (see Eqs. (2)–(5) and Table 1). Second special case takes place at  $\vartheta = 90^\circ$ ; formation of MDs becomes impossible, because there is no component of MDs Burgers vector contributing to the energy release during relaxation process [20–22,25]. These facts suggest that the use of substrates oriented in a certain way, for example, the use of  $\alpha$ - $\text{Al}_2\text{O}_3$  substrates with traditional (a-plane, m-plane) and exotic  $[(10\bar{1}7), (10\bar{1}6), (2\bar{1}\bar{1}11), (2\bar{1}\bar{1}9), \text{etc.}]$  growth planes, will make it possible to grow  $\kappa$ - $\text{Ga}_2\text{O}_3$  films with a low density of MDs.

As can be seen from Fig. 2, formation of pure edge MDs with  $b[100]$  and  $b[010]$  is energetically favorable (takes place at lower values of film thickness) than formation of mixed MDs with  $b[110]$ . However, in the case of heterostructures with inclination axis about  $[2\bar{1}\bar{1}0]$  crystallographic direction, the difference in thickness is insignificant (see Fig. 2a).

Due to similar crystal structure of  $\kappa$ - $\text{Al}_2\text{O}_3$  and  $\kappa$ - $\text{Ga}_2\text{O}_3$ , we analyze dependences of the critical thickness  $h_c$  on the inclination angle  $\vartheta$  for MDs formation in  $\kappa$ -( $\text{Al}_x\text{Ga}_{1-x}$ ) $_2\text{O}_3/\kappa$ - $\text{Al}_2\text{O}_3$  heterostructures with various Al composition  $x$ ; see Fig. 3. Solid, dashed and dash-dotted curves correspond to  $x = 0, 0.2,$  and  $0.4,$  respectively. Analyzing these dependences, one can conclude that an

increase in Al composition leads to an increase in the critical thickness for MDs formation. Despite the change in the Al composition, the formation of a pure edge MDs is energetically favorable than the formation of a mixed MDs in the considered heterostructures. Formation of pure edge MDs is easier in heterostructures with inclination axis about  $[010]$  crystallographic direction than in heterostructures with inclination axis about  $[100]$  crystallographic direction, and this is vice versa for mixed MDs.

#### 4. CONCLUSIONS

The analytical model that describes critical conditions for pure edge and mixed MDs formation in film/substrate heterostructures with various growth orientations has been modified for  $\kappa$ - $\text{Ga}_2\text{O}_3/\alpha$ - $\text{Al}_2\text{O}_3$  and  $\kappa$ -( $\text{Al}_x\text{Ga}_{1-x}$ ) $_2\text{O}_3/\kappa$ - $\text{Al}_2\text{O}_3$  heterostructures. Based on the symmetry of the crystal lattice of  $\kappa$ - $\text{Ga}_2\text{O}_3$ , we consider two possibilities for variation in growth orientations of the film in  $\kappa$ - $\text{Ga}_2\text{O}_3/\alpha$ - $\text{Al}_2\text{O}_3$  and  $\kappa$ - $\text{Ga}_2\text{O}_3/\kappa$ - $\text{Al}_2\text{O}_3$  heterostructures with inclination about either  $[100]$  or  $[010]$  crystallographic directions. Analytical expressions that describe the component of the lattice mismatch released by MDs have been proposed for heterostructures under consideration. It has been shown that for  $\kappa$ - $\text{Ga}_2\text{O}_3/\kappa$ - $\text{Al}_2\text{O}_3$  heterostructures there are two special orientations, for which the formation of MD is unfavorable: (i) standard orientation at  $\vartheta = 90^\circ$ , which is explained by the fact that there is no component of MDs Burgers vector contributing to the energy release during relaxation process; (ii) exotic orientation at  $\vartheta \sim 26^\circ$  (for heterostructure with inclination axis about  $[100]$  crystallographic direction of the film, or about  $[2\bar{1}\bar{1}0]$  crystallographic direction of the

substrate), and at  $\vartheta \sim 28^\circ$  (for heterostructure with inclination axis about [010] crystallographic direction of the film, or about  $[\bar{1}100]$  crystallographic direction of the substrate) that is explained by vanishing of the component of the lattice mismatch released by MDs due to features of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> crystal structure.

Calculations for  $\kappa$ -(Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub> heterostructures have demonstrated that an increase in Al composition  $x$  leads to an increase in the critical thickness for MDs formation. Formation of mixed MDs is easier in heterostructures with inclination axis about [100] crystallographic direction than in heterostructures with inclination axis about [010] crystallographic direction, and this is vice versa for pure edge MDs.

## ACKNOWLEDGEMENTS

A.V. Kremleva and A.Yu. Ivanov received support from Russian Science Foundation Project No. 21-79-00211.

## REFERENCES

- [1] S.I. Stepanov, V.I. Nikolaev, V.E. Bougrov, A.E. Romanov, *Gallium oxide: Properties and applications – A review*, Rev. Adv. Mater. Sci., 2016, vol. 44, no. 1, pp. 63–86.
- [2] A.A. Petrenko, Y.N. Kovach, D.A. Bauman, M.A. Odnoblyudov, V.E. Bougrov, A.E. Romanov, *Current state of Ga<sub>2</sub>O<sub>3</sub>-based electronic and optoelectronic devices. Brief review*, Rev. Adv. Mater. Technol., 2021, vol. 3, no. 2, pp. 1–26.
- [3] Y. Wang, J. Su, Z. Lin, J. Zhang, J. Chang, Y. Hao, *Recent progress on the effects of impurities and defects on the properties of Ga<sub>2</sub>O<sub>3</sub>*, J. Mater. Chem. C, 2022, Advance Article.
- [4] R. Roy, V.G. Hill, E.F. Osborn, *Polymorphism of Ga<sub>2</sub>O<sub>3</sub> and the System Ga<sub>2</sub>O<sub>3</sub>-H<sub>2</sub>O*, J. Am. Chem. Soc., 1952, vol. 74, no. 3, pp. 719–722.
- [5] S. Yoshioka, H. Hayashi, A. Kuwabara, F. Oba, K. Matsunaga, I. Tanaka, *Structures and energetics of Ga<sub>2</sub>O<sub>3</sub> polymorphs*, J. Phys.: Condens. Matter., 2007, vol. 19, no. 34, art. no. 346211.
- [6] H.Y. Playford, A.C. Hannon, E.R. Barney, R.I. Walton, *Structures of uncharacterised polymorphs of gallium oxide from total neutron diffraction*, Chem. Eur. J., 2013, vol. 19, no. 8, pp. 2803–2813.
- [7] Z. Galazka, R. Uecker, K. Irmscher, M. Albrecht, D. Klimm, M. Pietsch, M. Brützam, R. Bertram, S. Ganschow, R. Fornari, *Czochevski growth and characterization of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals*, Cryst. Res. Technol., 2010, vol. 45, no. 12, pp. 1229–1236.
- [8] A.V. Kremleva, S.S. Sharofidinov, A.M. Smirnov, E. Podlesnov, M.V. Dorogov, M.A. Odnoblyudov, V.E. Bougrov, A.E. Romanov, *Growth of thick gallium oxide on the various substrates by halide vapor phase epitaxy*, Mater. Phys. Mech., 2020, vol. 44, no. 2, pp. 164–171.
- [9] Y. Zhang, Y. Gong, X. Chen, Y. Kuang, J. Hao, F.-F. Ren, S. Gu, R. Zhang, J. Ye, *Unlocking the single-domain heteroepitaxy of orthorhombic  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> via phase engineering*, ACS Appl. Electron. Mater., 2022, vol. 4, no. 1, pp. 461–468.
- [10] I. Cora, F. Mezzadri, F. Boschi, M. Bosi, M. Čaplovičová, G. Calestani, I. Dódonny, B. Pécz, R. Fornari, *The real structure of  $\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub> and its relation to  $\kappa$ -phase*, CrystEngComm., 2017, vol. 19, no. 11, pp. 1509–1516.
- [11] H. Nishinaka, H. Komai, D. Tahara, Y. Arata, M. Yoshimoto, *Microstructures and rotational domains in orthorhombic  $\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub> thin films*, Jpn. J. Appl. Phys., 2018, vol. 57, no. 11, art. no. 115601.
- [12] Y.-H. Hsu, W.-Y. Wu, K.-L. Lin, Y.-H. Chen, Y.-H. Lin, P.-L. Liu, C.-L. Hsiao, R.-H. Horng,  *$\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub> grown on  $c$ -plane sapphire by MOCVD with a multistep growth process*, Cryst. Growth Des., 2022, vol. 22, no. 3, pp. 1837–1845.
- [13] K. Jiang, J. Tang, M.J. Cabral, A. Park, L. Gu, R.F. Davis, L.M. Porter, *Layered phase composition and microstructure of  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>-dominant heteroepitaxial films grown via MOCVD*, J. Appl. Phys., 2022, vol. 131, no. 5, art. no. 055305.
- [14] S. Shapenkov, O. Vyvenko, V. Nikolaev, S. Stepanov, A. Pechnikov, M. Scheglov, G. Varygin, *Polymorphism and faceting in Ga<sub>2</sub>O<sub>3</sub> layers grown by HVPE at various gallium-to-oxygen ratios*, Phys. Status Solidi B, 2022, vol. 259, no. 2, art. no. 2100331.
- [15] H. Nishinaka, O. Ueda, Y. Ito, N. Ikenaga, N. Hasuike, M. Yoshimoto, *Plan-view TEM observation of a single-domain  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> thin film grown on  $\varepsilon$ -GaFeO<sub>3</sub> substrate using GaCl<sub>3</sub> precursor by mist chemical vapor deposition*, Jpn. J. Appl. Phys., 2022, vol. 61, no. 1, art. no. 018002.
- [16] X. Hou, Y. Zou, M. Ding, Y. Qin, Z. Zhang, X. Ma, P. Tan, S. Yu, X. Zhou, X. Zhao, G. Xu, H. Sun, S. Long, *Review of polymorphous Ga<sub>2</sub>O<sub>3</sub> materials and their solar-blind photodetector applications*, J. Phys. D.: Appl. Phys., 2021, vol. 54, no. 4, art. no. 043001.
- [17] M. Kneiß, D. Splith, P. Schlupp, A. Hassa, H. von Wenckstern, M. Lorenz, M. Grundmann, *Realization of highly rectifying Schottky barrier diodes and pn heterojunctions on  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> by overcoming the conductivity anisotropy*, J. Appl. Phys., 2021, vol. 130, no. 8, art. no. 084502.
- [18] Y. Kuang, X. Chen, T. Ma, Q. Du, Y. Zhang, J. Hao, F.-F. Ren, B. Liu, S. Zhu, S. Gu, R. Zhang, Y. Zheng, J. Ye, *Band alignment and enhanced interfacial conductivity manipulated by polarization in a surfactant-mediated grown  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>O<sub>3</sub> heterostructure*, ACS Appl. Electron. Mater., 2021, vol. 3, no. 2, pp. 795–803.
- [19] S. Kim, J. Lee, H. Ahn, K. Kim, M. Yang, *Growth and crystal phase transformation of  $\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub> grown on 4H-SiC by MOCVD*, Mater. Sci. Semicond. Process., 2022, vol. 144, art. no. 106586.
- [20] A.E. Romanov, E.C. Young, F. Wu, A. Tyagi, C.S. Gallinat, S. Nakamura, S.P. DenBaars, J.S. Speck, *Basal plane misfit dislocations and stress relaxation in III-nitride semipolar heteroepitaxy*, J. Appl. Phys., 2011, vol. 109, no. 10, art. no. 103522.
- [21] A.M. Smirnov, E.C. Young, V.E. Bougrov, J.S. Speck, A.E. Romanov, *Critical thickness for the formation of misfit dislocations originating from prismatic slip in semipolar and nonpolar III-nitride heterostructures*, APL Mater., 2016, vol. 4, no. 1, art. no. 016105.
- [22] A.M. Smirnov, E.C. Young, V.E. Bougrov, J.S. Speck, A.E. Romanov, *Stress relaxation in semipolar and*



- nonpolar III-nitride heterostructures by formation of misfit dislocations of various origin*, J. Appl. Phys., 2019, vol. 126, no. 24, art. no. 245104.
- [23] A.M. Smirnov, A. V. Kremleva, S.S. Sharofidinov, V.E. Bougrov, A.E. Romanov, *Stress-strain state in  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> epitaxial films on  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> substrates*, Appl. Phys. Express, 2020, vol. 13, no. 7, art. no. 075502.
- [24] A.M. Smirnov, A.V. Kremleva, Sh.Sh. Sharofidinov, V.E. Bougrov, A.E. Romanov, *Misfit stress relaxation in  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> heterostructures via formation of misfit dislocations*, Phys. Solid State., 2021, vol. 63, no. 6, pp. 915–922.
- [25] A.M. Smirnov, A.V. Kremleva, Sh.Sh. Sharofidinov, A.E. Romanov, *Misfit stress relaxation in wide bandgap semiconductor heterostructures with trigonal and hexagonal crystal structure*, J. Appl. Phys., 2022, vol. 131, no. 2, art. no. 025301.
- [26] L.B. Freund, S. Suresh, *Thin Film Materials: Stress, Defect Formation and Surface Evolution*, Cambridge University Press, Cambridge, 2008.
- [27] J.W. Matthews, A.E. Blakeslee, *Defects in epitaxial multilayers*, J. Cryst. Growth., 1974, vol. 27, pp. 118–125.
- [28] L.B. Freund, *Dislocation Mechanisms of Relaxation in Strained Epitaxial Films*, MRS Bull., 1992, vol. 17, no. 7, pp. 52–60.
- [29] K. Shimada, *First-principles study of crystal structure, elastic stiffness constants, piezoelectric constants, and spontaneous polarization of orthorhombic Pna2<sub>1</sub>-M<sub>2</sub>O<sub>3</sub> (M = Al, Ga, In, Sc, Y)*, Mater. Res. Express., 2018, vol. 5, no. 3, art. no. 036502.
- [30] D.B. Hovis, A. Reddy, A.H. Heuer, *X-ray elastic constants for  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>*, Appl. Phys. Lett., 2006, vol. 88, no. 13, art. no. 131910.

УДК 548.4

## Релаксация напряжений за счет образования дислокаций в орторомбических пленках Ga<sub>2</sub>O<sub>3</sub>, полученных на подложках Al<sub>2</sub>O<sub>3</sub>

А.М. Смирнов<sup>1</sup>, А. Ю. Иванов<sup>1</sup>, А.В. Кремлева<sup>1</sup>, Ш.Ш. Шарофидинов<sup>2</sup>, А.Е. Романов<sup>1,3</sup>

<sup>1</sup> Институт перспективных систем передачи данных Университета ИТМО, Кронверкский пр., д. 49, лит А, Санкт-Петербург 197101, Россия

<sup>2</sup> Сектор твердотельной электроники ФТИ им. Иоффе РАН, ул. Политехническая, д. 26, Санкт-Петербург, 194021, Россия

<sup>3</sup> Сектор теории твердого тела ФТИ им. Иоффе РАН, ул. Политехническая, д. 26, Санкт-Петербург, 194021, Россия

**Аннотация.** Проанализирована предпочтительность образования различных типов дислокаций несоответствия (ДН) в гетероструктурах типа пленка/подложка  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> и  $\kappa$ -(Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub>. Рассмотрены две возможные ориентации роста пленок, определяемые углом наклона  $\vartheta$ , в упомянутых выше гетероструктурах с осями наклона относительно кристаллографических направлений [100] или [010]. Исследованы зависимости критической толщины пленки от угла наклона  $\vartheta$  при зарождении ДН в рассматриваемых гетероструктурах. Установлено наличие двух особых ориентаций ( $\vartheta \sim 26^\circ$  для гетероструктуры [100],  $\vartheta \sim 28^\circ$  для гетероструктуры [010] и  $\vartheta = 90^\circ$  для обоих типов наклона) гетероструктур  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, при которых образование ДН энергетически невыгодно. Показано, что в гетероструктурах [010]  $\kappa$ -(Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub>/ $\kappa$ -Al<sub>2</sub>O<sub>3</sub> образование краевых ДН предпочтительнее, чем в гетероструктурах [100], а для смешанных ДН – наоборот.

**Ключевые слова:** гетероструктуры типа пленка/подложка; оксид галлия; ориентация роста; дислокации несоответствия